

ODES: LECTURES 11+12

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What is an ODE/System of ODEs?

- ODEs are equation(s) for continuous function(s) of a single independent variable involving derivatives of the unknown function(s).

Normally written

$$y' = f(t, y) \quad ' = \frac{d}{dt}$$

$$\begin{bmatrix} \frac{dy_1}{dt} \\ \vdots \\ \frac{dy_n}{dt} \end{bmatrix} = \begin{bmatrix} f_1(t, y_1, y_2, \dots, y_n) \\ \vdots \\ f_n(t, y_1, y_2, \dots, y_n) \end{bmatrix}$$

NOTE:

~~Higher order ODEs~~ Higher order ODEs can always be transformed into equivalent first order systems

EX $\frac{d^2 y}{dx^2} = y t$ Put $y' = y_1$ $\Rightarrow \frac{dy_1}{dx} = y_1 t$
 $\frac{dy_1}{dx} = \frac{dy}{dx} = y_2$ $\frac{dy_1}{dx} = y_2$

$$\Rightarrow \begin{bmatrix} y_1' \\ y_2' \end{bmatrix} = \begin{bmatrix} y_2 \\ y_1 t \end{bmatrix}$$

An ODE/Sys. of ODEs does not determine a unique solution (it only specifies y'). Usually, there are infinitely many solns.

To single out a particular soln we also have to specify a value, i.e. give (a) boundary condition(s).

The type of boundary conditions needed/desired define the problem as either:

- 1) An initial value problem (IVP)
- or 2) A boundary value problem (BVP)

EX (IVP) $y' = -y$ with $y(t_0) = y_0$
 $y'' = y^3 t$ with $y(t_0) = y_0$ $y'(t_0) = d_0$

EX (BVP) $y'' = y^3 t$ with $y(t_0) = y_0$ $y(t_1) = y_1$
 $y''' = -y^2$ with $y(t_0) = y_0$ $y'(t_1) = d_1$

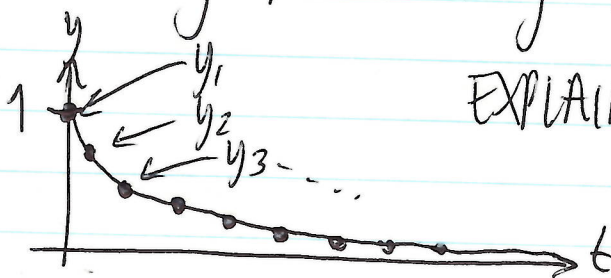
Initial Value problems

(3)

$$y' = f(t, y) \text{ with } y(t_0) = y_0$$

A numerical soln. to this problem consists of finding approximate values of the soln at discrete points in time.

EX $y' = -y$ with $y(0) = 1 \Rightarrow y = \exp(-t)$.



EXPLAIN: h_k, t_k notation.

The central question is: how should we accurately and quickly compute these approximate values?

A: Discretise the derivative operator!

Different discretisations of ~~the~~ $\frac{d}{dt}$ give rise to different time-stepping schemes with different properties.

Stable and unstable ODEs

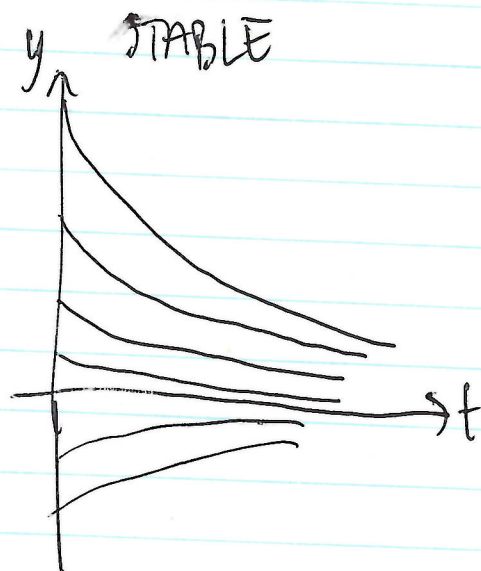
If the members of a family move away from one another as t increases $\rightarrow$ unstable

———— " ———— closer ———— $\rightarrow$ stable.

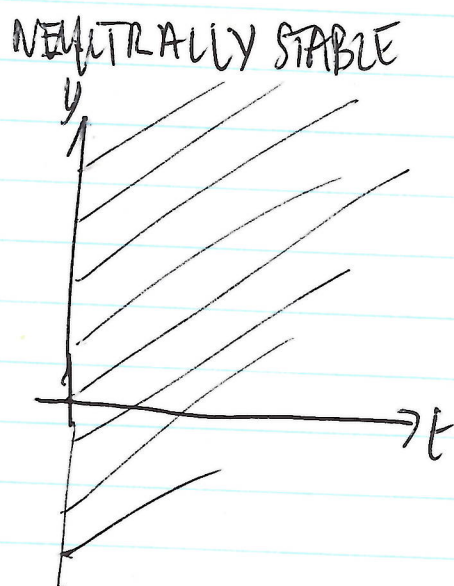
If neither $\rightarrow$ neutrally stable.

More precisely the system $y' = f(t, y)$ is unstable if $J_{ij} = \frac{\partial f_i}{\partial x_j}$ has any eigenvalues > 0
 stable if ———— " ———— has only ^{eigenvalues} < 0

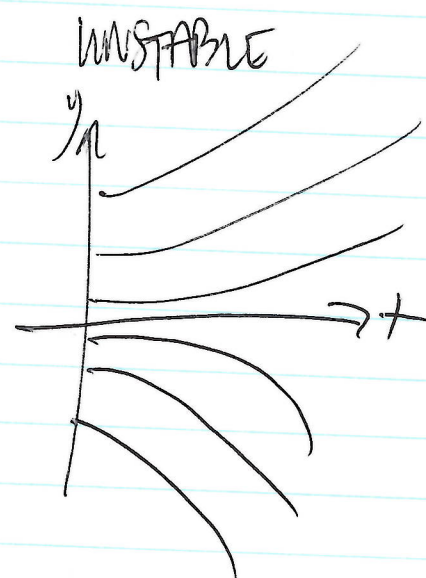
EX $y' = -y$



$y' = 1$



$y' = y$



Euler's Method

$y(t+h)$ - is the object we wish to find given $y(t)$.

Expand $y(t+h)$ as a Taylor Series about $y(t) \Rightarrow$

$$y(t+h) = y(t) + y'(t)h + \frac{y''(t)}{2}h^2 + O(h^3).$$

If we truncate at $O(h^2)$ we have the following approximation:

$$y'(t) \approx \frac{y(t+h) - y(t)}{h} \quad \text{~~not~~}$$

In the discrete setting

$$y_{i+1} = y_i + h_i f(t_i, y_i).$$

NOTE: In general $y_i \neq y|_{t=t_i}$ $\leftarrow$ The method is approximate.

Accuracy and Stability

5

Like all other numerical procedures ~~different~~ time stepping methods suffer from 2 sources of error

1) Rounding

2) Truncation.

We can't do much about 1), 2) we can!

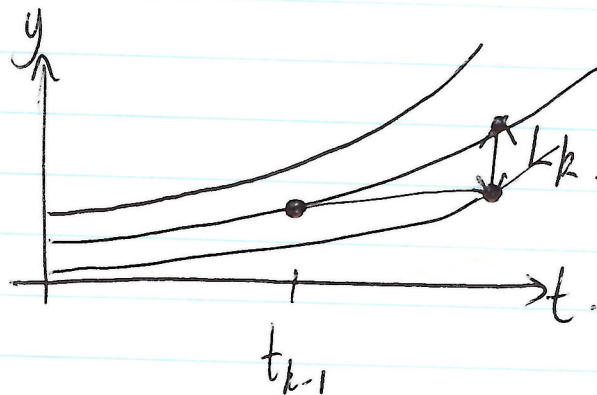
Truncation errors

(a) Local truncation error. (The error made in making one time step).

$$L_k = y_k - u_{k-1}(t_k).$$

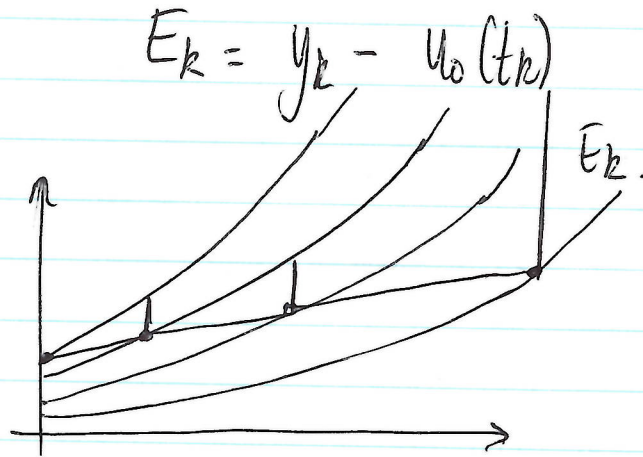
The computed soln at time t_k .

(exact)
The real solution curve that passes through (y_{k-1}, t_{k-1}) .

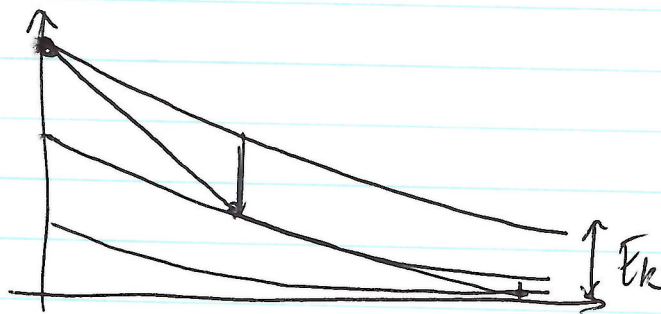


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(b) Global truncation error - Difference between computed soln. and true soln. determined by initial data.



NOTE: $E_k \neq \sum_{i=1}^k L_i$



However a small E_k is what we really are about. However all we can control is L_k .

Accuracy: A scheme is said to be of order p if $L_k = O(h_k^{p+1})$.

~~more and more~~

stability of a num. meth.

Different from the concept of stability of the ODE being solved!

A method is stable if small perturbations lead to small changes in the computed solns.

Unstable if small ~~change~~ perturbation leads to big changes.

~~One~~ One way to assess stability is to look at the model equation $y' = \lambda y$ with $y(0) = y_0$

$$\Rightarrow y = y_0 \exp(-\lambda t).$$

EX: Stability for forward Euler.

If we choose a fixed step-size h , Euler forward says

$$\begin{aligned} y_{k+1} &= y_k + y_k \lambda h \\ &= y_k (1 + \lambda h). \end{aligned}$$

Thus

$$y_k = (1 + \lambda h)^k y_0$$

The exact soln $\rightarrow 0$ as $t \rightarrow \infty$. As does the numerical method, provided $\lambda < 0$. ~~that forward Euler is stable if also~~

(9)

This says that the errors do not grow provided

$$|1 + \lambda h| < 1 \Rightarrow \lambda h \in (0, -2)$$

$$\lambda h < -2 \\ h > -2/\lambda$$

In this interval FEuler stable
otherwise unstable.

Unstable if $\lambda > 0 \leftarrow$ the equation itself is unstable

Unstable if $h > \frac{2}{\lambda} \leftarrow$ the equation is stable, but the time step too large. The integration is unstable.

Backward Euler

$$y_{k+1} = y_k + f(t_{k+1}, y_{k+1})h_k$$

Very similar to FEuler, but one important difference. ~~we need the value of $f(t_{k+1}, y_{k+1})$ before we know y_{k+1}~~

Can no longer simply evaluate the RHS. In fact, this is a system of equations to solve for y_{k+1} .

This makes Backward E. implicit whereas Forward E. explicit.