

Q1

a) Probability $P = \frac{V_C}{V_B} = \frac{\text{Volume of the circles}}{\text{Volume of the box}}$

Volume of the box
(independent of n): $V_B = L^2$

Volume of the circle for $n=1$ $V_{C(1)} = \pi R^2 = \pi \left(\frac{L}{2}\right)^2 = \frac{\pi L^2}{4}$

————— 11 ————— for $n=2$

$V_{C(2)} = \underbrace{2^2}_{\substack{\# \text{ of} \\ \text{circles}}} \pi \left(\frac{L}{2 \cdot 2}\right)^2 = 4 \pi \frac{L^2}{16} = \frac{\pi L^2}{4}$

Volume of the circle for $n=4$

$V_{C(4)} = \underbrace{4^2}_{\substack{\# \text{ of} \\ \text{circles}}} \pi \left(\frac{L}{2 \cdot 4}\right)^2 = 16 \pi \frac{L^2}{64} = \frac{\pi L^2}{4}$

and for an arbitrary n

$V_{C(n)} = n^2 \cdot \pi \left(\frac{L}{2 \cdot n}\right)^2 = \cancel{n^2} \pi \frac{L^2}{4 \cdot \cancel{n^2}} = \underline{\underline{\frac{\pi L^2}{4}}}$

Therefore, the converging probability

$P_n = \frac{V_{C(n)}}{V_B} = \frac{\pi \frac{L^2}{4}}{L^2} = \frac{\pi}{4} = \lim_{n \rightarrow \infty} P_n$

(independent of n !)

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b) Now a sphere in a 3D box

$$\text{Volume of a sphere: } V_s = \frac{4}{3}\pi R^3$$

$$\text{Volume of a 3D box: } V_B = L^3$$

Radius of a sphere
(when n spheres are packed
in every direction)

$$R = \frac{L/2}{n} = \frac{L}{2n}$$

of spheres in a 3D box: n^3

Volume of spheres

$$V_{S(n)} = n^3 \cdot \frac{4}{3}\pi \left(\frac{L}{n \cdot 2}\right)^3 = \cancel{n^3} \frac{4}{3}\pi \frac{L^3}{\cancel{n^3} 8} = \frac{\pi}{6} L^3$$

(again, independent
of n)

$$P = \frac{V_{S(n)}}{V_B} = \frac{\frac{\pi}{6} L^3}{L^3} = \frac{\pi}{6}$$

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c) Volume of a k -dimensional sphere

$$V_k = \frac{\pi^{\frac{k}{2}}}{(\frac{k}{2})!} R^k$$

Note that this is valid only for k even
(for k odd the expression for V is different)

Note that for $k=2$

$$V_2 = \frac{\pi}{1} R^2 = \pi R^2 \quad (\text{see a})$$

Radius of a k -sphere

(when n spheres packed
in every direction)

$$R = \frac{L/2}{n} = \frac{L}{2 \cdot n}$$

of spheres in a k -dimensional box: n^k

(HW1/2)

Volume of the spheres:

$$V_{\text{sph}} = n^k \cdot \frac{\pi^{\frac{k}{2}}}{(\frac{k}{2})!} \left(\frac{L}{2 \cdot n}\right)^k = \cancel{n^k} \frac{\pi^{\frac{k}{2}}}{(\frac{k}{2})!} \frac{L^k}{2^k \cdot \cancel{n^k}}$$

again, independent of n

Volume of a k -dimensional box:

$$V_B = \frac{L^k}{\pi^{k/2}} \frac{L^k}{2^k}$$

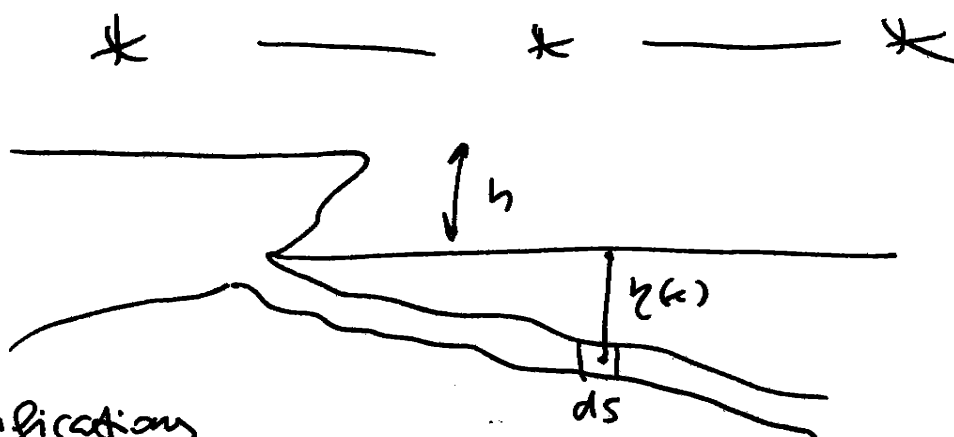
Probability:

$$P = \frac{V_{\text{sph}}}{V_B} = \frac{\frac{\pi^{k/2}}{(\frac{k}{2})!} \frac{L^k}{2^k}}{L^k} =$$

Note that the result of point (a) is obtained for $k=2$

$$= \frac{\pi^{k/2}}{(\frac{k}{2})! 2^k}$$

Q2



Modifications

- potential energy of the fluid must be taken into account

~~the fluid must be taken into account~~

Potential energy:

New energy $E = E_K + E_P$, $E_K = \frac{1}{2} \rho A x(\dot{x})^2$ (olol)

Potential energy of an infinitesimal element ds

$$dE_P = \rho g [h(x) - h] A ds$$

where h is an arbitrary reference level

~~the fluid must be taken into account~~

Then

$$E_p = \int_0^{x(t)} dE_p = \int_0^{x(t)} \rho g [\eta(x') - h] A dx' = \rho g A \int_0^{x(t)} (\eta(x') - h) dx'$$

Now the time derivative:

$$\frac{d}{dt} (E_p + E_k) = \frac{d}{dt} E_p + \frac{d}{dt} E_k$$

$\frac{dE_k}{dt}$ was calculated at the lecture and is given by

$$\frac{dE_k}{dt} = \rho A \dot{x} \left(\frac{1}{2} \dot{x}^2 + x \ddot{x} \right)$$

Then using the chain rule

$$\frac{dE_p}{dt} = \frac{d}{dx} \left[\rho g A \int_0^x (\eta(x') - h) dx' \right] \frac{dx}{dt} = \rho g A [\eta(x) - h] \dot{x}$$

By the
Fundamental
Theorem of
Calculus

Hence, finally

$$\frac{dE}{dt} = \rho A \dot{x} \left[\frac{1}{2} \dot{x}^2 + x \ddot{x} + g(\eta(x) - h) \right]$$

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Solve the equation using separation of variables

$$\frac{dx}{dt} = x^2 \Rightarrow \int x^{-2} dx = \int dt$$

$$-\frac{1}{x} = t + C \Rightarrow x(t) = -\frac{1}{t + C}$$

Using the initial condition $x(0) = x_0$

$$x(0) = -\frac{1}{C} = x_0 \Rightarrow C = -\frac{1}{x_0}$$

Hence: $x(t) = -\frac{1}{t - \frac{1}{x_0}}$

We note that

$$\lim_{t \rightarrow \left(\frac{1}{x_0}\right)^+} x(t) = \infty$$

Therefore the solution $x(t)$ is defined only for
 $t \in [0, \frac{1}{x_0})$ and ceases to exist at $t = \underline{\underline{\frac{1}{x_0}}}$