

1) Use the index notation and Einstein summation convention

a) start with the terms with the curl

$$\begin{aligned} [\bar{a} \times (\nabla \times \bar{b})]_i &= \epsilon_{ijk} a_j \epsilon_{kpq} \frac{\partial b_q}{\partial x_p} \\ &= \epsilon_{ijk} \epsilon_{kpq} a_j \frac{\partial b_q}{\partial x_p} = \epsilon_{kij} \epsilon_{kpq} a_j \frac{\partial b_q}{\partial x_p} = \\ &\quad \epsilon_{ijk} = \epsilon_{kij} \end{aligned}$$

Now use the following property of the Levi-Civita tensor

$$\epsilon_{kij} \epsilon_{kpq} = \delta_{ip} \delta_{jq} - \delta_{iq} \delta_{jp}$$

$$= (\delta_{ip} \delta_{jq} - \delta_{iq} \delta_{jp}) a_j \frac{\partial b_q}{\partial x_p} = a_q \frac{\partial b_q}{\partial x_i} - a_p \frac{\partial b_i}{\partial x_p}$$

change $\left. \begin{array}{l} q \rightarrow p \\ \text{in the} \\ \text{first} \\ \text{term} \end{array} \right\} \rightarrow$

$$= a_p \frac{\partial b_p}{\partial x_i} - a_p \frac{\partial b_i}{\partial x_p}$$

Likewise:

$$[\bar{b} \times (\nabla \times \bar{a})]_i = b_p \frac{\partial a_p}{\partial x_i} - b_p \frac{\partial a_i}{\partial x_p}$$

The other two terms on the RHS

$$[(\vec{a} \cdot \nabla) \vec{b}]_i = a_p \frac{\partial b_i}{\partial x_p}, \quad [(\vec{b} \cdot \nabla) \vec{a}]_i = b_p \frac{\partial a_i}{\partial x_p}$$

Putting together all the terms on the RHS

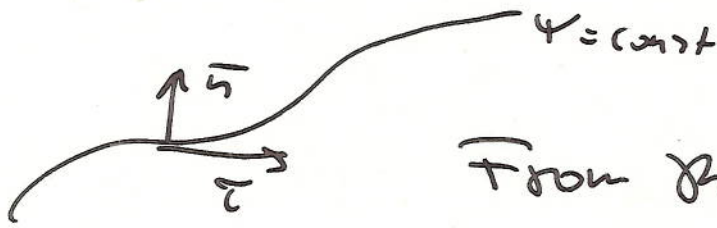
$$\begin{aligned} [\vec{a} \times (\nabla \times \vec{b}) + \vec{b} \times (\nabla \times \vec{a}) + (\vec{a} \cdot \nabla) \vec{b} + (\vec{b} \cdot \nabla) \vec{a}]_i &= \\ &= a_p \frac{\partial b_p}{\partial x_i} - \cancel{a_p \frac{\partial b_i}{\partial x_p}} + b_p \frac{\partial a_p}{\partial x_i} - \cancel{b_p \frac{\partial a_i}{\partial x_p}} \\ &\quad + \cancel{a_p \frac{\partial b_i}{\partial x_p}} + \cancel{b_p \frac{\partial a_i}{\partial x_p}} = \frac{\partial}{\partial x_i} (a_p b_p) \\ &= [\nabla (\vec{a} \cdot \vec{b})]_i \end{aligned}$$



$$\begin{aligned} \text{b) } \left[\frac{1}{2} (\nabla \times \vec{a}) \times \vec{b} \right]_i &= \frac{1}{2} \epsilon_{pqr} \frac{\partial a_m}{\partial x_q} \epsilon_{ipk} b_k \\ &= -\frac{1}{2} \epsilon_{pqr} \epsilon_{pik} \frac{\partial a_m}{\partial x_q} b_k = -\frac{1}{2} (\delta_{qi} \delta_{mk} - \delta_{qk} \delta_{mi}) \frac{\partial a_m}{\partial x_q} b_k \\ &= -\frac{1}{2} \left(\frac{\partial a_k}{\partial x_i} - \frac{\partial a_i}{\partial x_k} \right) b_k = \underbrace{\frac{1}{2} \left(\frac{\partial a_i}{\partial x_k} - \frac{\partial a_k}{\partial x_i} \right)}_{R_{ik}} b_k \\ &= [\vec{R} \vec{b}]_i \end{aligned}$$



2) WLOG consider 2D case and introduce a local coordinate system attached to the streamlines



From the momentum equation

$$\nabla \left(\underbrace{\frac{1}{\rho} p + \frac{1}{2} |\bar{u}|^2 + \phi}_{H = H(n)} \right) = \bar{u} \times (\nabla \times \bar{u})$$

$$\nabla = \bar{e}_s \frac{\partial}{\partial s} + \bar{e}_n \frac{\partial}{\partial n}$$

~~$$\nabla H = \bar{e}_s \frac{\partial H}{\partial s} + \bar{e}_n \frac{\partial H}{\partial n} = \bar{e}_n \frac{\partial H}{\partial n}$$~~

~~$$\nabla H = \bar{e}_s \frac{\partial H}{\partial s} + \bar{e}_n \frac{\partial H}{\partial n} = \bar{e}_n \frac{\partial H}{\partial n}$$~~

$$\bar{u} \times (\nabla \times \bar{u}) = \left[\bar{u} \cdot (\bar{u} \times (\nabla \times \bar{u})) \right] \bar{u} + 0 \bar{e}_s$$

$$\nabla H = \underbrace{\bar{e}_s \frac{\partial H}{\partial s}}_0 + \bar{e}_n \frac{\partial H}{\partial n} = \left[\bar{u} \cdot (\bar{u} \times (\nabla \times \bar{u})) \right] \bar{u}$$

Thus:

$$\frac{\partial H}{\partial n} = \bar{u} \cdot (\bar{u} \times (\nabla \times \bar{u}))$$

derivative in the direction normal to the stream-function

4.4) [Solution to problem 3 on the last page]

a) $\vec{u} = \nabla \phi$ in Ω

Equation satisfied by the potential:

$$\nabla \cdot \vec{u} = 0 \Rightarrow \Delta \phi = 0 \text{ in } \Omega$$

BCs: $\frac{\partial \phi}{\partial n} = 0$ at $y=0, b$

$$\frac{\partial \phi}{\partial n} = -\frac{\partial \phi}{\partial x} = -V \text{ at } x=0$$

$$\frac{\partial \phi}{\partial n} = \frac{\partial \phi}{\partial x} = V \text{ at } x=a$$

} (*)

The solution of Neumann problem (*) is ~~unique~~

$$\phi(x, y) = Vx + \phi_0, \quad \phi_0 - \text{arbitrary constant}$$

The velocity field

$$\vec{u} = \nabla \phi = \underline{[V, 0]}$$

b) $\vec{u} = \nabla \perp \psi$ in Ω

$$\frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} = 0 \Rightarrow \nabla \perp \nabla \perp \psi = \Delta \psi = 0$$

BCs

at ~~boundary~~

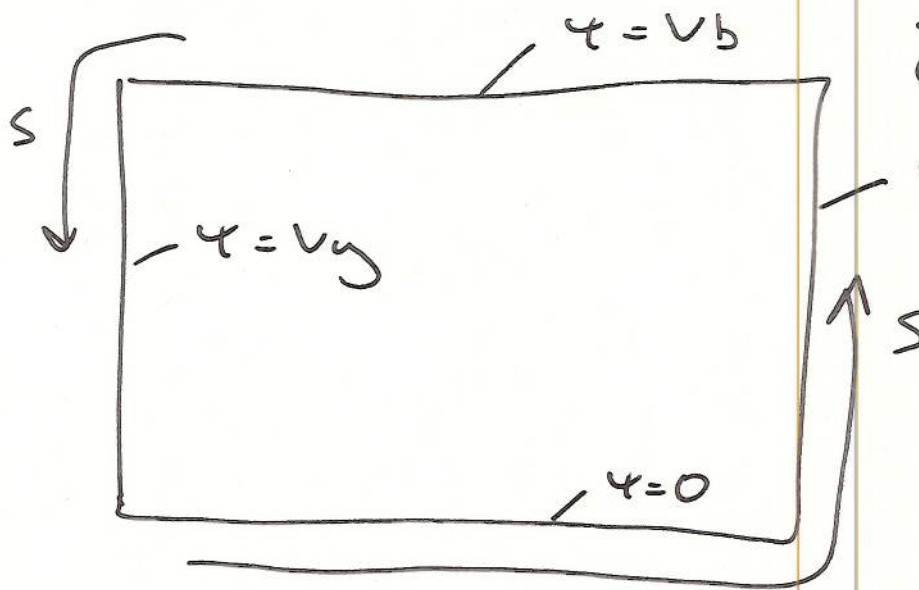
$$y=0 \quad \nabla \perp \psi \cdot \vec{n} = \frac{\partial \psi}{\partial s} = 0 \Rightarrow \psi = \psi_0 = 0$$

(arbitrary constant)

$$x=a \quad \nabla \perp \psi \cdot \vec{n} = \frac{\partial \psi}{\partial s} = V$$

$$\psi(a, y) = \int_0^y V ds = Vy + \psi_0 = Vy$$

(4)



Boundary conditions for stream function ψ

$$y=b \quad \nabla \perp \psi \cdot \vec{n} = \frac{\partial \psi}{\partial s} = 0 \Rightarrow \psi = \text{const} = Vb$$

$$x=0 \quad \nabla \perp \psi \cdot \vec{n} = \frac{\partial \psi}{\partial s} = -V$$

$$\psi(y) = -\int_0^{b-y} V ds' = \left| \begin{array}{l} s' = b-y' \\ ds' = -dy' \end{array} \right|$$

$$= V \int_b^y dy' = V(y-b) + \underbrace{Vb}_{\psi_0} = Vy$$

↑ ensures continuity

Thus, ψ satisfies the following Dirichlet problem

$$\begin{cases} \Delta \psi = 0 & \text{in } \Omega \\ \psi = Vy & \text{at } x=0, a \\ \psi = 0 & \text{at } y=0 \\ \psi = Vb & \text{at } y=b \end{cases}$$

The solution is $\psi(x,y) = Vy$

$$\text{So that } \vec{u} = \nabla \perp \psi = \left[\frac{\partial \psi}{\partial y}, -\frac{\partial \psi}{\partial x} \right] = [V, 0]$$

(5)

3) ~~Streamlines~~ Equation for the streamlines

$$\frac{dx}{u} = \frac{dy}{v}$$

a)
$$\frac{dx}{-y/\sqrt{x^2+y^2}} = \frac{dy}{x/\sqrt{x^2+y^2}} \Rightarrow xdx = -ydy$$

$$\frac{1}{2}x^2 + \frac{1}{2}y^2 = C \Rightarrow x^2 + y^2 = C'$$

Or streamlines are circles

b)
$$\frac{dx}{-y} = \frac{dy}{x} \Rightarrow xdx = -ydy \Rightarrow x^2 + y^2 = C'$$

Or streamlines are also circles

Field (a) is irrotational ($\omega = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$)

$$\frac{\partial}{\partial x} \left(\frac{x}{x^2+y^2} \right) - \frac{\partial}{\partial y} \left(\frac{-y}{x^2+y^2} \right) = 0$$

Field (b) has nonvanishing vorticity

$$\frac{\partial}{\partial x} (x) - \frac{\partial}{\partial y} (-y) = 2 \neq 0$$