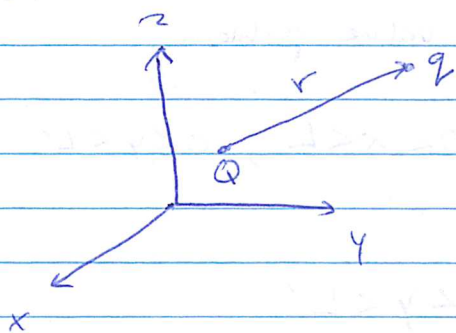


## Electrostatic Potential



Suppose we have two points with <sup>positive</sup> charges  $Q$  and  $q$  which are  $r$  units apart.

• Coulomb's law states that each repels the other with force whose magnitude is  $F = \frac{qQ}{4\pi\epsilon_0 r^2}$ , where  $\epsilon_0$  = permittivity of free space (it's a constant)

• We can define the vector field (called electric field intensity) by

$$\mathbf{E} = \frac{Q}{4\pi\epsilon_0 r^3} \mathbf{r}$$

where  $\mathbf{r}$  is the vector from  $Q$  to a unit charge  $q$ .

• We showed in examples from chapter 14 that  $\nabla \times \mathbf{E} = \mathbf{0} = \nabla \cdot \mathbf{E}$ .

•  $\nabla \times \mathbf{E} = \mathbf{0} \Rightarrow$  on a simply connected domain,  $\mathbf{E} = -\nabla V$  for some potential function  $V$ .

• Since  $\nabla \cdot \mathbf{E} = 0 \Rightarrow \nabla \cdot (-\nabla V) = -\nabla^2 V = 0 \Rightarrow \nabla^2 V = 0$  which is Laplace's Equation.

• For this specific  $\mathbf{E}$ , we can simply solve  $\mathbf{E} = -\nabla V$  to find  $V$ .

For more general charge distributions, Laplace's equation still holds outside the distribution, and Poisson's equation ( $\nabla^2 V = \frac{-\rho}{\epsilon_0}$ ) holds inside. More details are provided in section 20.3

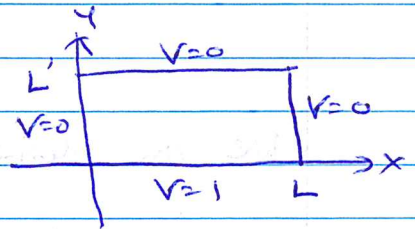
• You have seen Laplace's equation in Math 2130. Solutions to this differential equation were said to be harmonic.

Example: Solve the boundary value problem

$$\Delta V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0, \quad 0 < x < L, \quad 0 < y < L',$$

$$\begin{aligned} V(0, y) = 0 = V(L, y) & \quad 0 < y < L' \\ V(x, L') = 0, \quad V(x, 0) = 1 & \quad 0 < x < L \end{aligned}$$

for potential in the rectangular plate when sides are maintained at the above potentials.



Using separation of variables, write  $V(x, y) = X(x)Y(y)$  so that

$$\frac{d^2 X}{dx^2} Y + \frac{d^2 Y}{dy^2} X = 0$$

$$\Rightarrow \frac{X''}{X} = -\frac{Y''}{Y}$$

$$\Rightarrow X'' + \lambda X = 0 \quad \text{and} \quad Y'' - \lambda Y = 0 \quad \text{using } -\lambda \text{ in the separation principle.}$$

The boundary conditions imply  $X(0) = 0 = X(L)$  and  $Y(L') = 0$ .

Using information from our computations in chapter 19, we see that  $\lambda_n = \frac{n^2 \pi^2}{L^2}$ ,  $n \geq 1$  are eigenvalues for

$X'' + \lambda X = 0$  and a solution has the form

$$C \sin\left(\frac{n\pi x}{L}\right) \quad \text{for some } n \geq 1 \text{ integer}$$

The second system  $Y'' - \lambda Y = 0$  has  
 auxiliary equation  $m^2 - \lambda = 0$ , but  $\lambda$  from  
 the previous computation is  $\frac{n^2 \pi^2}{L^2}$ , so  $m = \pm \frac{n\pi}{L}$

Using our knowledge of ODE solutions, we see that a general  
 solution is

$$Y(y) = D e^{\frac{n\pi y}{L}} + F e^{-\frac{n\pi y}{L}}$$

The boundary condition  $\Rightarrow F = -D e^{\frac{2n\pi L}{L}}$

$$\Rightarrow Y(y) = D \left[ e^{\frac{n\pi y}{L}} - e^{\frac{n\pi(2L-y)}{L}} \right]$$

$$\Rightarrow V(x,y) = X(x)Y(y) \\ = b \sinh\left(\frac{n\pi x}{L}\right) \left[ e^{\frac{n\pi y}{L}} - e^{\frac{n\pi(2L-y)}{L}} \right]$$

with  $b \in \mathbb{R}$ . This is a solution to Laplace's equation, and  
 since all of the conditions are homogeneous, we can use  
 superposition to get

$$V(x,y) = \sum_{n=1}^{\infty} b_n \sinh\left(\frac{n\pi x}{L}\right) \left[ e^{\frac{n\pi y}{L}} - e^{\frac{n\pi(2L-y)}{L}} \right]$$

The final condition  $V(x,0) = 1$  implies

$$\sum_{n=1}^{\infty} C_n \sinh\left(\frac{n\pi x}{L}\right) = 1, \quad 0 < x < L$$

with  $C_n = b_n (1 - e^{\frac{2n\pi L}{L}})$ .

• We can recognize this as the eigenfunction expansion of  
 the function 1.

$$\text{Then } C_n = \frac{2}{L} \int_0^L (1) \sinh\left(\frac{n\pi x}{L}\right) dx = \frac{2 \left[ 1 + (-1)^{n+1} \right]}{n\pi}$$

$$\Rightarrow b_n = \frac{2 \left[ 1 + (-1)^{n+1} \right]}{(1 - e^{\frac{2n\pi L}{L}})}$$

This series can be simplified (?) to the one on  
line (21.26).

Try working through Part C in 21.2 for practice  
with potential problems.