Open Questions in Quantum Information Geometry

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1. Classical Parametric Information Geometry

- Study of differential geometric properties of families of classical probability densities.

- Given a probability space (Ω, Σ, μ), a family of probability densities $\mathcal{M} = \{p(x, \theta)\}$, for sample points $x \in \Omega$ and parameters $\theta = (\theta^1, \ldots, \theta^n) \in \mathbb{R}^n$ can be viewed as a Riemannian manifold equipped with the Fisher metric

$$g_{ij} = \int \frac{\partial \log p(x,\theta)}{\partial \theta^i} \frac{\partial \log p(x,\theta)}{\partial \theta^j} p(x,\theta) dx \tag{1}$$

Apart from the Levi-Civita connection associated with g, the statistical manifold $\mathcal{M}$ can be equipped with the exponential connection

$$\left(\nabla^{(1)}_{\frac{\partial}{\partial\theta^i}} \frac{\partial}{\partial\theta^j}\right)(p) = \frac{\partial^2 \log p}{\partial\theta^i \partial\theta^j} - E_p\left(\frac{\partial^2 \log p}{\partial\theta^i \partial\theta^j}\right),$$

and the mixture connection

$$\left(\nabla^{(-1)}_{\frac{\partial}{\partial\theta^i}} \frac{\partial}{\partial\theta^j}\right)(p) = \frac{\partial^2 \log p}{\partial\theta^i \partial\theta^j} + \frac{\partial \log p}{\partial\theta^i}\frac{\partial \log p}{\partial\theta^j},$$

which are dual to the metric g in the sense that $\langle\cdot,\cdot\rangle_p$ if

$$v\left(g(s_1, s_2)\right) = g(\nabla^{(1)}_v s_1, s_2) + g(s_1, \nabla^{(-1)}_v s_2) \tag{2}$$

for all $v \in T_p\mathcal{M}$ and all smooth vector fields s_1 and s_2.

One can also define a family of α–connections induced by the embeddings

$$\begin{aligned} \ell_\alpha \ &: \ \mathcal{M} \to \mathcal{A} \\ & p \mapsto \frac{2}{1-\alpha} p^{\frac{1-\alpha}{2}}, \end{aligned}$$

where $\mathcal{A}$ is the algebra of random variables on Ω and prove that they satisfy

$$\nabla^\alpha = \frac{1+\alpha}{2}\nabla^{(1)} + \frac{1-\alpha}{2}\nabla^{(-1)}. \tag{3}$$

The Fisher metric is the unique Riemannian metric reduced by all stochastic maps on the tangent bundle to $\mathcal{M}$. The duality of the α–connections with respect to it lead to rich minimization/projection theorems related to their associated α–divergences.

2. Classical Nonparametric Information Geometry

2.1. Classical Orlicz Spaces

Consider Young functions of the form

$$\Phi(x) = \int_0^{|x|} \phi(t)dt, \quad x \geq 0, \tag{4}$$

where $\phi : [0, \infty) \mapsto [0, \infty)$ is nondecreasing, continuous and such that $\phi(0) = 0$ and $\lim_{x \to \infty} \phi(x) = +\infty$. This include the monomials $|x|^r/r$, for $1 < r < \infty$, as well as the following examples:

$$\begin{aligned} \Phi_1(x) &= \cosh x - 1, & (5) \\ \Phi_2(x) &= e^{|x|} - |x| - 1, & (6) \\ \Phi_3(x) &= (1 + |x|) \log(1 + |x|) - |x| & (7) \end{aligned}$$

The complementary of a Young function Φ of the form (4) is given by

$$\Psi(y) = \int_0^{|y|} \psi(t)dt, \quad y \geq 0, \tag{8}$$

where ψ is the inverse of ϕ. One can verify that (Φ_2, Φ_3) and $(|x|^r/r, |x|^s/s)$, with $r^{-1} + s^{-1} = 1$, are examples of complementary pairs.

Now let (Ω, Σ, P) be a probability space. The Orlicz space associated with a Young function Φ defined as

$$L^{\Phi}(P) = \left\{ f : \Omega \mapsto \overline{\mathbf{R}}, \text{measurable} : \int_\Omega \Phi(\alpha f) dP < \infty, \text{ for some } \alpha > 0 \right\}.$$

If we identify functions which differ only on sets of measure zero, then L^{Φ} is a Banach space when furnished with the Luxembourg norm

$$N_{\Phi}(f) = \inf\left\{k > 0 : \int_{\Omega} \Phi(\frac{f}{k})dP \leq 1\right\}, \tag{9}$$

or with the equivalent Orlicz norm

$$\|f\|_{\Phi} = \sup\left\{\int_{\Omega} |fg|d\mu : g \in L^{\Psi}(\mu), \int_{\Omega} \Psi(g)dP \leq 1\right\}, \tag{10}$$

where Ψ is the complementary Young function to Ψ.

2.2. The Pistone-Sempi Manifold

Consider the set

$$\mathcal{M} \equiv \mathcal{M}(\Omega, \Sigma, \mu) = \{f : \Omega \mapsto \mathbf{R}, f > 0 \text{ a.e. and } \int_\Omega f d\mu = 1\}.$$

For each point $p \in \mathcal{M}$, let $L^{\Phi_1}(p)$ be the exponential Orlicz space over the probability space $(\Sigma, \Omega, p d\mu)$ and consider its closed subspace of p-centred random variables

$$B_p = \{u \in L^{\Phi_1}(p) : \int_\Omega u p d\mu = 0\} \tag{11}$$

as the coordinate Banach space.

In probabilistic terms, the set $L^{\Phi_1}(p)$ correspond to random variables whose moment generating function with respect to the probability $p d\mu$ is finite on a neighborhood of the origin.

They define one dimensional exponential models $p(t)$ associated with a point $p \in \mathcal{M}$ and a random variable u:

$$p(t) = \frac{e^{tu}}{Z_p(tu)} p, \qquad t \in (-\varepsilon, \varepsilon). \tag{12}$$

Define the inverse of a local chart around $p \in \mathcal{M}$ as

$$\begin{aligned} e_p : \mathcal{V}_p &\rightarrow \mathcal{M} \\ u &\mapsto \frac{e^u}{Z_p(u)} p. \end{aligned} \tag{13}$$

Denote by $\mathcal{U}_p$ the image of $\mathcal{V}_p$ under e_p. Let e_p^{-1} be the inverse of e_p on $\mathcal{U}_p$. Then a local chart around p is given by

$$\begin{aligned} e_p^{-1} : \mathcal{U}_p &\rightarrow B_p \\ q &\mapsto \log\left(\frac{q}{p}\right) - \int_\Omega \log\left(\frac{q}{p}\right) p d\mu. \end{aligned} \tag{14}$$

For any $p_1, p_2 \in \mathcal{M}$, the transition functions are given by

$$\begin{aligned} e_{p_2}^{-1} e_{p_1} : e_{p_1}^{-1}(\mathcal{U}_{p_1} \cap \mathcal{U}_{p_2}) &\rightarrow e_{p_2}^{-1}(\mathcal{U}_{p_1} \cap \mathcal{U}_{p_2}) \\ u &\mapsto u + \log\left(\frac{p_1}{p_2}\right) - \int_\Omega \left(u + \log\frac{p_1}{p_2}\right) p_2 d\mu. \end{aligned}$$

Proposition 1 *For any $p_1, p_2 \in \mathcal{M}$, the set $e_{p_1}^{-1}(\mathcal{U}_{p_1} \cap \mathcal{U}_{p_2})$ is open in the topology of B_{p_1}.*

We then have that the collection $\{(\mathcal{U}_p, e_p^{-1}), p \in \mathcal{M}\}$ satisfies the three axioms for being a C^∞–atlas for $\mathcal{M}$. Moreover, since all the spaces B_p are isomorphic as topological vector spaces, we can say that $\mathcal{M}$ is a C^∞–manifold modeled on $B_p \equiv T_p\mathcal{M}$.

Given a point $p \in \mathcal{M}$, the connected component of $\mathcal{M}$ containing p coincides with the maximal exponential model obtained from p: $\mathcal{E}(p) = \left\{\frac{e^u}{Z_p(u)} p, u \in B_p\right\}$.

2.3. The Fisher Information and Dual Connections

Let $\langle \cdot, \cdot \rangle_p$ be a continuous positive definite symmetric bilinear form assigned continuously to each $B_p \simeq T_p\mathcal{M}$. A pair of connection (∇, ∇^*) are said to be dual with respect to $\langle \cdot, \cdot \rangle_p$ if

$$\langle \tau u, \tau^* v \rangle_q = \langle u, v \rangle_p \tag{15}$$

for all $u, v \in T_p\mathcal{M}$, where τ and τ^* denote the parallel transports associated with ∇ and ∇^*, respectively.

Equivalently, (∇, ∇^*) are dual with respect to $\langle \cdot, \cdot \rangle_p$ if

$$v\left(\langle s_1, s_2 \rangle_p\right) = \langle \nabla_v s_1, s_2 \rangle_p + \langle s_1, \nabla^*_v s_2 \rangle_p \tag{16}$$

for all $v \in T_p\mathcal{M}$ and all smooth vector fields s_1 and s_2.

The infinite dimensional generalisation of the Fisher information is given by

$$\langle u, v\rangle_p = \int_\Omega (uv)p d\mu, \qquad \forall u, v \in B_p. \tag{17}$$

This is clearly bilinear, symmetric and positive definite. Moreover, continuity follows from that fact that, since $L^{\Phi_1}(p) \simeq L^{\Phi_2}(p) \subset L^{\Phi_3}(p)$, the generalised Hölder inequality gives

$$|\langle u, v\rangle_p| \leq K\|u\|_{\Phi_1,p}\|v\|_{\Phi_1,p}, \qquad \forall u, v \in B_p. \tag{18}$$

If p and q are two points on the same connected component of $\mathcal{M}$, then the exponential parallel transport is given by

$$\begin{aligned} \tau_{pq}^{(1)} : T_p\mathcal{M} &\rightarrow T_q\mathcal{M} \\ u &\mapsto u - \int_\Omega uq\,d\mu. \end{aligned} \tag{19}$$

To obtain duality with respect to the Fisher information, we define the mixture parallel transport on $T\mathcal{M}$ as

$$\begin{aligned} \tau_{pq}^{(-1)} : T_p\mathcal{M} &\rightarrow T_q\mathcal{M} \\ u &\mapsto \frac{p}{q}u, \end{aligned} \tag{20}$$

for p and q in the same connected component of $\mathcal{M}$.

Theorem 2 *The connections $\nabla^{(1)}$ and $\nabla^{(-1)}$ are dual with respect to the Fisher information.*

2.4. α–connections

We begin with Amari's α-embeddings

$$\begin{aligned} \ell_\alpha : \mathcal{M} &\rightarrow L^r(\mu) \\ p &\mapsto \frac{2}{1-\alpha} p^{\frac{1-\alpha}{2}}, \quad \alpha \in (-1,1), \end{aligned} \tag{21}$$

where $r = \frac{2}{1-\alpha}$. Observe that $\ell_\alpha(p) \in S^r(\mu)$, the sphere of radius r in $L^r(\mu)$.

We are now ready to define the α-connections. In what follows, $\widetilde{\nabla}$ is used to denote the trivial connection on $L^r(\mu)$.

Definition 3 *For $\alpha \in (-1,1)$, let $\gamma : (-\varepsilon, \varepsilon) \to \mathcal{M}$ be a smooth curve such that $p = \gamma(0)$ and $v = \dot{\gamma}(0)$ and let $s \in S(T\mathcal{M})$ be a differentiable vector field. The α-connection on $T\mathcal{M}$ is given by*

$$\left(\nabla_v^\alpha s\right)(p) = (\ell_\alpha)_{*(p)}^{-1}\left[\Pi_{rp^{1/r}}\widetilde{\nabla}_{(\ell_\alpha)_{*(p)}v}(\ell_\alpha)_{*(\gamma(t))}s\right]. \tag{22}$$

Theorem 4 *The exponential, mixture and α-covariant derivatives on $T\mathcal{M}$ satisfy*

$$\nabla^\alpha = \frac{1+\alpha}{2}\nabla^{(1)} + \frac{1-\alpha}{2}\nabla^{(-1)}. \tag{23}$$

Corollary 5 *The connections ∇^α and $\nabla^{-\alpha}$ are dual with respect to the Fisher information $\langle\cdot,\cdot\rangle_p$.*

3. Finite Dimensional Quantum Systems

- $\mathcal{H}^N$: finite dimensional complex Hilbert space;

- $\mathcal{B}(\mathcal{H}^N)$: algebra of operators on $\mathcal{H}^N$;

- $\mathcal{A}$: N^2-dimensional real vector subspace of self-adjoint operators;

- $\mathcal{M}$: n-dimensional submanifold of all invertible density operators on $\mathcal{H}^N$, with $n = N^2 - 1$.

3.1 Quantum α-connections

For $\alpha \in (-1, 1)$, define the α-embedding of $\mathcal{M}$ into $\mathcal{A}$ as

$$\begin{aligned} \ell_\alpha \ &: \ \mathcal{M} \to \mathcal{A} \\ & \rho \mapsto \frac{2}{1-\alpha}\rho^{\frac{1-\alpha}{2}}. \end{aligned}$$

At each point $\rho \in \mathcal{M}$, consider the subspace of $\mathcal{A}$ defined by

$$\mathcal{A}_\rho^{(\alpha)} = \left\{ A \in \mathcal{A} : \mathrm{Tr}\left(\rho^{\frac{1+\alpha}{2}} A\right) = 0 \right\},$$

and define the isomorphism

$$\begin{aligned} (\ell_\alpha)_{*(\rho)} \ &: \ T_\rho\mathcal{M} \to \mathcal{A}_\rho^{(\alpha)} \\ & v \mapsto (\ell_\alpha \circ \gamma)'(0). \end{aligned} \qquad (24)$$

Let $r = \frac{2}{1-\alpha}$. If we equip $\mathcal{A}$ with the the r-norm

$$\|A\|_r := (\mathrm{Tr}|A|^r)^{1/r},$$

then the α-embedding can be vied as a mapping from $\mathcal{M}$ to the positive part of the sphere of radius r.

Definition 6 *For $\alpha \in (-1, 1)$, let $\gamma : (-\varepsilon, \varepsilon) \to \mathcal{M}$ be a smooth curve such that $\rho = \gamma(0)$ and $v = \dot{\gamma}(0)$ and let $s \in S(T\mathcal{M})$ be a differentiable vector field. The α-connection on $T\mathcal{M}$ is given by*

$$\left(\nabla_v^{(\alpha)} s\right)(\rho) = (\ell_\alpha)_{*(\rho)}^{-1} \left[\Pi_{r\rho^{1/r}} \widetilde{\nabla}_{(\ell_\alpha)_{*(\rho)} v} (\ell_\alpha)_{*(\gamma(t))} s\right], \qquad (25)$$

where $\Pi_{r\rho^{1/r}}$ is the canonical projection from the tangent space $T_{r\rho^{1/r}}\mathcal{A} = \mathcal{A}$ onto the tangent space $T_{r\rho^{1/r}} S^r = \mathcal{A}^{(\alpha)}$.

3.2. Monotone Metrics

Now let us consider the extended manifold of faithful weights $\widehat{\mathcal{M}}$ (the positive definite matrices) and use the -1-representation (the limiting case $\alpha = -1$ of the α-representations) in order to define a Riemannian metric g on $\widehat{\mathcal{M}}$ by means of the inner product $\langle \cdot, \cdot \rangle_\rho$ in $\mathcal{A} \subset B(\mathcal{H}^N)$. We say that $\widehat{g}$ is monotone if and only if

$$\left\langle S(A^{(-1)}), S(A^{(-1)}) \right\rangle_{S(\rho)} \leq \left\langle A^{(-1)}, A^{(-1)} \right\rangle_\rho , \qquad (26)$$

for every $\rho \in \mathcal{M}$, $A \in T_\rho\mathcal{M}$, and every completely positive, trace preserving map $S : \mathcal{A} \to \mathcal{A}$.

For any metric $\widehat{g}$ on $T\widehat{\mathcal{M}}$, define the positive (super) operator K_σ on $\mathcal{A}$ by

$$\widehat{g}_\rho(\widehat{A}, \widehat{B}) = \left\langle \widehat{A}^{(-1)}, K_\rho\left(\widehat{B}^{(-1)}\right)\right\rangle_{HS} = \mathrm{Tr}\left(\widehat{A}^{(-1)} K_\rho\left(\widehat{B}^{(-1)}\right)\right). \quad (27)$$

Define also the (super) operators, $L_\rho X := \rho X$ and $R_\rho X := X\rho$, for $X \in \mathcal{A}$, which are also positive.

Theorem 7 (Petz 96) *A Riemannian metric g on $\mathcal{A}$ is monotone if and only if*

$$K_\sigma = \left(R_\sigma^{1/2} f(L_\sigma R_\sigma^{-1}) R_\sigma^{1/2}\right)^{-1},$$

where K_σ is defined in (27) and $f : R^+ \to R^+$ is an operator monotone function satisfying $f(t) = t f(t^{-1})$.

In particular, the BKM (Bogoluibov–Kubo–Mori) metric

$$g_\rho^B(A,B) = \int_0^\infty \mathrm{Tr}\left(\frac{1}{t+\rho}A^{(-1)}\frac{1}{t+\rho}B^{(-1)}\right)dt \tag{28}$$

and the WYD (Wigner-Yanase-Dyson) metric

$$g_\rho^{(\alpha)}(A,B) := \mathrm{Tr}\left(A^{(\alpha)}B^{(-\alpha)}\right), \qquad A,B \in T_\rho\mathcal{M}, \tag{29}$$

for $\alpha \in (-1,1)$ are special cases of monotone metrics corresponding respectively to the operator monotone functions

$$f^B(t) = \frac{t-1}{\log t}$$

and

$$f_p(x) = \frac{p(1-p)(x-1)^2}{(x^p-1)(x^{1-p}-1)}$$

for $p = \frac{1+\alpha}{2}$.

Theorem 8 *If the connections* $\nabla^{(1)}$ *and* $\nabla^{(-1)}$ *are dual with respect to a monotone Riemannian metric* g *on* $\mathcal{M}$, *then* g *is a* scalar *multiple of the BKM metric.*

Theorem 9 *If the connections* $\nabla^{(\alpha)}$ *and* $\nabla^{(-\alpha)}$ *are dual with respect to a monotone Riemannian metric* $\widehat{g}$ *on* $\widehat{\mathcal{M}}$, *then* $\widehat{g}$ *is a* scalar *multiple of the* WYD *metric.*

Corollary 10

$$\nabla^{\alpha} \neq \frac{1+\alpha}{2}\nabla^{(1)} + \frac{1-\alpha}{2}\nabla^{(-1)}. \tag{30}$$

3.3 Scalar curvature

- It has been proposed that the scalar curvature associated with a monotone metric represents the "average uncertainty" for a quantum state. In this way, it should be an increasing function under stochastic maps.

- The Bures metric is known to be a counter-example for $n \geq 3$ (Ditmann 99), while being zero for $n = 2$.

- Counter-examples are now known even for $n = 2$ (Andai 03).

- Petz's conjecture: The scalar curvature for the BKM metric is monotone (convincing numerical and theoretical evidence, but no proof !)

- Gibilisco and Isola's conjecture: The scalar curvature for the WYD metric for α close to -1 is monotone (motivated by α–geometry).

- Research proposal: Look at the scalar curvature associated with dual connections, instead of the Levi-Civita one.

4. Infinite dimensional quantum systems

- $\mathcal{H}$: infinite dimensional complex Hilbert space;
- $\mathcal{B}(\mathcal{H})$: algebra of operators on $\mathcal{H}$;
- $\mathcal{C}_p, 0 < p < 1$: compact operators $A : \mathcal{H} \mapsto \mathcal{H}$ such that $|A|^p \in \mathcal{C}_1$, where $\mathcal{C}_1$ is the set of trace-class operators on $\mathcal{H}$. Define

$$\mathcal{C}_{<1} := \bigcup_{0<p<1} \mathcal{C}_p.$$

- $\mathcal{M} = \mathcal{C}_{<1} \cap \Sigma$ where $\Sigma \subset \mathcal{C}_1$ denotes the set of normal faithful states on H.

4.1. ε-Bounded Perturbations

Let $H_0 \geq I$ be a self-adjoint operator with domain $\mathcal{D}(H_0)$, quadratic form q_0 and form domain $Q_0 = \mathcal{D}(H_0^{1/2})$, and let $R_0 = H_0^{-1}$ be its resolvent at the origin.

For $\varepsilon \in (0, 1/2)$, let $\mathcal{T}_\varepsilon(0)$ be the set of all symmetric forms X defined on Q_0 and such that $\|X\|_\varepsilon(0) := \left\| R_0^{\frac{1}{2}+\varepsilon} X R_0^{\frac{1}{2}-\varepsilon} \right\|$ is finite.

Then the map $A \mapsto H_0^{\frac{1}{2}-\varepsilon} A H_0^{\frac{1}{2}+\varepsilon}$ is an isometry from the set of all bounded self-adjoint operators on $\mathcal{H}$ onto $\mathcal{T}_\varepsilon(0)$. Hence $\mathcal{T}_\varepsilon(0)$ is a Banach space with the ε-norm $\|\cdot\|_\varepsilon(0)$.

Lemma 11 *For fixed symmetric X, $\|X\|_\varepsilon$ is a monotonically increasing function of $\varepsilon \in [0, 1/2]$.*

4.2. Construction of the Manifold

To each $\rho_0 \in \mathcal{C}_{\beta_0} \cap \Sigma$, $\beta_0 < 1$, let $H_0 = -\log \rho_0 + cI \geq I$ be a self-adjoint operator with domain $\mathcal{D}(H_0)$ such that

$$\rho_0 = Z_0^{-1} e^{-H_0} = e^{-(H_0 + \Psi_0)}. \tag{31}$$

In $\mathcal{T}_\varepsilon(0)$, take X such that $\|X\|_\varepsilon(0) < 1 - \beta_0$. Since $\|X\|_0(0) \leq \|X\|_\varepsilon(0) < 1 - \beta_0$, X is also q_0-bounded with bound a_0 less than $1 - \beta_0$. The *KLMN* theorem then tells us that there exists a unique semi-bounded self-adjoint operator H_X with form $q_X = q_0 + X$ and form domain $Q_X = Q_0$. Following an unavoidable abuse of notation, we write $H_X = H_0 + X$ and consider the operator

$$\rho_X = Z_X^{-1} e^{-(H_0 + X)} = e^{-(H_0 + X + \Psi_X)}. \tag{32}$$

Then $\rho_X \in \mathcal{C}_{\beta_X} \cap \Sigma$, where $\beta_X = \frac{\beta_0}{1-a_0} < 1$ [Streater 2000]. We take as a neighbourhood $\mathcal{M}_0$ of ρ_0 the set of all such states, that is, $\mathcal{M}_0 = \{\rho_X : \|X\|_\varepsilon(0) < 1 - \beta_0\}$.

The map

$$\rho_X \mapsto \widehat{X} = X - \rho_0 \cdot X$$

is then a global chart for the Banach manifold $\mathcal{M}_0$ modeled by $\widehat{\mathcal{T}}_\varepsilon(0) = \{X \in \mathcal{T}_\varepsilon(0) : \rho_0 \cdot X = 0\}$. We extend our manifold by adding new patches compatible with $\mathcal{M}_0$.

4.3. Affine Structure and Analyticity of the Free Energy

Given two points ρ_X and ρ_Y in $\mathcal{M}_0$ and their tangent spaces $\widehat{\mathcal{T}}_\varepsilon(X)$ and $\widehat{\mathcal{T}}_\varepsilon(Y)$, we define the torsion free, flat, $(+1)$-parallel transport along any continuous path γ connecting ρ_X and ρ_Y in the manifold as

$$\tau^(1) : Z \in \widehat{\mathcal{T}}_\varepsilon(X) \mapsto (Z - \rho_Y \cdot Z) \in \widehat{\mathcal{T}}_\varepsilon(Y)$$

Open problem: How to define the (-1) parallel transport ?

Theorem 12 *The free energy of the state* $\rho_X = Z_X^{-1} e^{-H_X} \in \mathcal{C}_{\beta_X} \subset \mathcal{M}, \beta_X < 1$*, defined by*

$$\Psi(\rho_X) := \log Z_X, \tag{33}$$

is infinitely Fréchet differentiable and has a convergent Taylor series for sufficiently small neighbourhoods of ρ_X *in* $\mathcal{M}$*.*

As a consequence, we can define the infinite dimensional BKM metric as

$$g_\rho(X, Y) = \int_0^1 \mathrm{Tr}\left(X \rho^\lambda Y \rho^{(1-\lambda)}\right) d\lambda \tag{34}$$

5. Noncommutative Orlicz Spaces

- $\mathcal{A}$: semifinite von Neumann algebra of operators on $\mathcal{H}$ with a faithful semifinite normal trace τ.

- $L^0(\mathcal{A}, \tau)$: closed densely defined operators $x = u|x|$ affiliated with $\mathcal{A}$ with the property that, for each $\varepsilon > 0$, there exits a projection $p \in \mathcal{A}$ such that $p\mathcal{H} \subset \mathcal{D}(x)$ and $\tau(1 - p) \leq \varepsilon$ (trace measurable operators).

- $\tilde{x}(t) := \inf\{s > 0 : \tau(e_{(s,\infty)}) \leq t\}$, where $e_{(\cdot)}$ are the spectral projections of $|x|$ (rearrangement function).

Lemma 13 (Fack/Kosaki) *Let* $0 \leq a \in L^0(\mathcal{A}, \tau)$*. Then*

$$\tau(\phi(a)) = \int_0^\infty \phi(\tilde{a}(t))dt$$

for any continuous increasing function $\phi : [0, \infty) \to [0, \infty)$*.*

Definition 14 (Kunze) *The Orlicz space associated with* $(\mathcal{A}, \tau, \phi)$ *is*

$$L^\phi(\mathcal{A}, \tau) = \{x \in L^0(\mathcal{A}, \tau) : \tau(\phi(\lambda|x|)) \leq 1, \text{ for some } \lambda > 0\}.$$

For Orlicz spaces associated with a state, one has

Definition 15 (Zegarlinski) *Given a state* $\omega(a) = \tau(\rho a)$, *for an invertible density operator* ρ, *the Orlicz space associated with* $(\mathcal{A}, \omega, \phi)$ *is*

$$L^{\phi}(\mathcal{A}, \tau) = \{x \in L^{0}(\mathcal{A}, \tau) : O(\lambda x) \leq 1, \text{ for some } \lambda > 0\},$$

where, for a given $s \in [0, 1]$,

$$O(x) = \tau(\phi(|(\phi^{-1}(\rho))^{s} x (\phi^{-1}(\rho))^{1-s}|)).$$

Question: What is the classical analogue for $\phi = cosh(x) - 1$ and $\mathcal{A}$ commutative ?

Alternatively,

Definition 16 (Streater) *A map* Φ *from* H_0*–bounded forms to* $R_+ \cup \infty$ *is called a quantum Young function if it is convex, even, satisfy* $\Phi(0) = 0$ *and* $\Phi(X) > 0$ *if* $X \neq 0$*, and is finite for all* X *with* sufficiently *small Kato bounds.*

Lemma 17 *The function*

$$\Phi(X) = \frac{1}{2} Tr\left(e^{-H_0-\Psi_0-X} + e^{-H_0-\Psi_0+X}\right) - 1$$

is a quantum Young function.

One can then define a quantum analogue of the Luxemburg norm and prove analogue's of Young's and Hölder's inequality (using the BKM metric).

Question: Are the norms on overlapping charts equivalent ?